

RBS-study of $\text{Ge}_x\text{Si}_{1-x}$ Compounds Formed by Variable Dose Ge Implantation into Si Wafers

A. Ramírez^a, A. Zehe^{b*}, A. Thomas^c

^a*Benemérita Universidad Autónoma de Puebla, Instituto de Ciencias,
17 oriente # 1603, Puebla, Pue., México*

^b*Facultad de Cs. Físico-Matemáticas, Apdo.*

Post. # 1505, 72000 Puebla, Pue., México

^c*Disetronic GmbH, Sulzbach/Ts, Germany*

Received: November 12, 2001; Revised: April 14, 2002

Amorphous and relaxed epitaxial GeSi films are prepared by Ge-implantation into Si(111) wafers of both 60 keV and 200 keV energetic Ge⁺-ions with appropriate dose, followed by post-implantation thermal annealing, comprising a single final annealing at a temperature of 900 °C. The implantation dose was varied between 10¹⁴ and 10¹⁷ atoms cm⁻². Rutherford backscattering (RBS) and channeling analysis was applied in order to explore the formation of a single crystalline Si-Ge compound layer, both prior and after the thermal treatment. The depth and the thickness of the implanted layer, as well as their molar composition and crystalline quality was determined, and it was found that a single crystalline Si-Ge alloy layer was created, with both depth and mole fraction depending on the ion energy and the ion dose.

Keywords: *Germanium implantation, solid phase epitaxy, Rutherford backscattering, siGe/si heterostructure*

1. Introduction

GeSi/Si heterostructure devices play an important role in microelectronics concepts since the heterojunction fabrication technology has shown noticeable improvements. Unfortunately, difficulties with the incorporation of MBE or CVD based GeSi growth techniques into current very large scale integration schemes have shown to delay the commercialization of GeSi devices. An alternative approach for fabricating GeSi/Si heterostructures may be a technique based on high dose germanium ion implantation into silicon, since ion implantation is a well established VLSI technique characterized by a high wafer throughput. Up to now, ion implantation has been employed extensively in the additive microfabrication process of electronic devices, where the chemical, physical or electrical properties of semiconductors are altered by bombarding the substrate with high energetic ions of the appropriate chemical species. Adjustment of the ion energy allows to modify the range of the ions before they come to rest, permitting control of the depth distribution of the added material. Research on GeSi/Si

heterostructure materials has led recently to dramatic improvements in performance and functionality of silicon-based electronic and optoelectronic devices¹⁻⁵.

Germanium and silicon are miscible over the whole composition range, thus $\text{Ge}_x\text{Si}_{1-x}$ is an ideal material to study alloy effects at any value of the mole fraction x . Over the past several years, advances in the Si/Ge technology have brought about an impressive potential for fabricating high-performance Si/Si_{1-x}Ge_x heterojunction bipolar transistors⁶⁻⁸ as well as other Si/SiGe heterojunction devices⁹. However, even though Si/SiGe heterojunction devices can outperform bulk Si devices in a variety of aspects, the higher complexity and cost of Si/Ge epitaxial growth have yet prevented laboratory technologies from entering into high volume production¹⁰. Techniques like molecular beam epitaxy (MBE)^{11,12}, low-pressure chemical vapor deposition (LP CVD) and ultrahigh vacuum chemical vapor deposition (UHV CVD)¹³⁻¹⁵, atmospheric pressure chemical vapor deposition (AP CVD)¹⁶ and very low pressure chemical vapor deposition (VLP CVD)¹⁷ are intensely studied and have been demonstrated to deliver high quality heteroepitaxial materials in

*e-mail: azehe@prodigy.net.mx

an extremely cost competitive silicon semiconductor industry. Alternative methods for the production of single-crystalline Ge-Si films were also developed. Abelson *et al.*¹⁸ produced epitaxial $\text{Ge}_x\text{Si}_{1-x}/\text{Si}(100)$ structures by evaporating Ge on Si(100) substrates and a subsequent laser beam mixing. The composition and the thickness of the alloy layers depend on the deposited Ge film thickness and the laser power. A similar process is established by the irradiation of the Ge-Si system with an electron beam¹⁹. First results on the damage induced by Ge^+ -ion implantation into Si were reported by Mersey *et al.*²⁰. The possibility of producing epitaxial $\text{Ge}_x\text{Si}_{1-x}/\text{Si}$ structures by heavy-dose Ge^+ -implantation and subsequent laser annealing was investigated in several occasions²¹⁻²³. Multiple germanium implantation provided also for a higher Ge incorporation²⁶.

The aim of the present work is, to deal with the formation of a single crystalline heterosystem $\text{Ge}_x\text{Si}_{1-x}/\text{Si}(111)$ by Ge^+ -ion implantation into silicon wafers in dependence on ion energy and ion dose, respectively. Rutherford backscattering (RBS) and channeling technique is applied prior and after a thermal treatment in order to study the crystalline character of the system.

2. Experiments

(111) oriented Si wafers, n-doped to reach a resistivity of $8 \Omega\text{cm}$, fixed with a small misorientation of about two degrees on a massive metal base without additional cooling were implanted with Ge^+ -ions at room temperature. The pressure in the ion accelerator was 1×10^{-4} Pa. Ions of two different energies were generated, 60 keV and 200 keV, respectively. The ion dose was varied between 1×10^{14} and $1 \times 10^{17} \text{ cm}^{-2}$ and applied while the beam was in a scanning mode. Thereafter the samples were annealed at $900 \text{ }^\circ\text{C}$ for 60 minutes in a nitrogen atmosphere. Prior to and after the annealing process the samples were studied with Rutherford backscattering spectroscopy (RBS), using $10.7 \text{ MeV } ^4\text{He}^+$ -ions in both the random and the channeling (aligned) crystal directions. The incident charge of the probing beam was $10 \mu\text{As}$, and the typical beam spot diameter was kept at 1 mm. The thickness and the depth of the implanted layer as well as their composition and crystalline quality were determined.

3. Results and Discussion

The distribution of the implanted germanium ions was calculated by the Monte-Carlo simulation program TRIM (transport of ions in matter)²⁴ for three different ion energies, 60 keV, 200 keV, and 300 keV, respectively, and is shown in Figs. 1 (a) and (b). The RBS spectra for the 'as implanted' samples with 60 and 200 keV Ge^+ -ion energies, are shown in the Figs. 2 and 3, respectively.

The random spectra are displayed only for the lowest

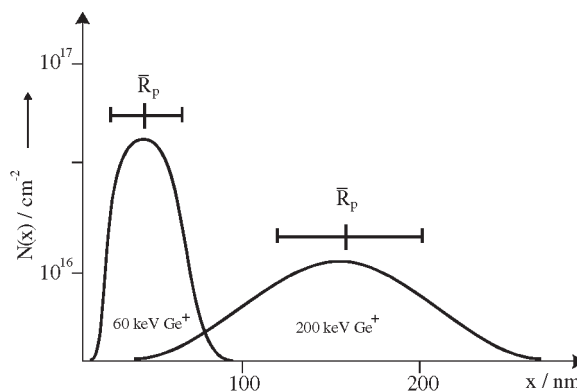


Figure 1a. Distribution in depth $N(x)$ of implanted Ge^+ -ions into Si, calculated by the simulation program TRIM²⁴ with a step width of $\Delta x = 10 \text{ nm}$. An implantation dose of 10^{17} cm^{-2} was chosen in the calculus.

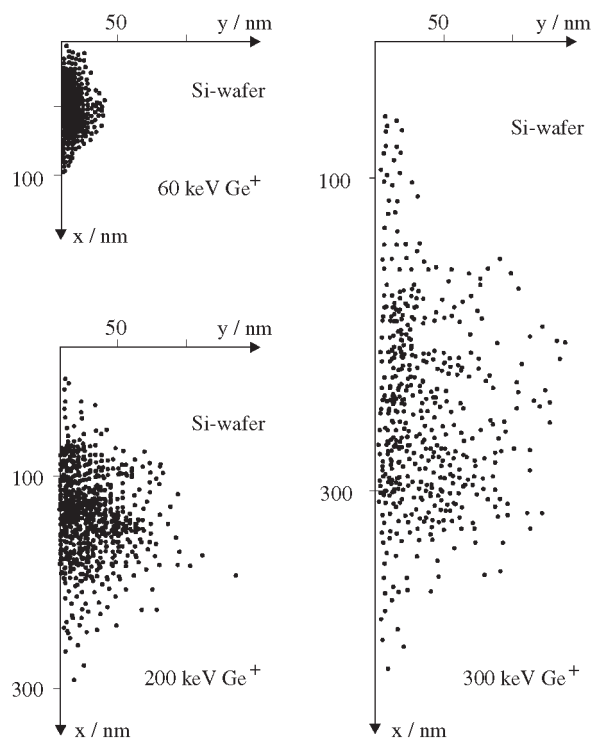


Figure 1b. Final position x of 500 implanted germanium ions as calculated by the simulation program TRIM²⁴, applying ion energies of 60 keV, 200 keV, and 300 keV. The lateral y -distribution results from a projection of the spatial ion distribution onto the x - y plane.

dose of $1 \times 10^{14} \text{ cm}^{-2}$ being representative for all doses up to $1 \times 10^{17} \text{ cm}^{-2}$. Considering first the implantation energy of 60 keV, it is apparent for the aligned spectra that even the lowest dose causes substantial radiation damage, as can be judged from the peak occurring between channel numbers

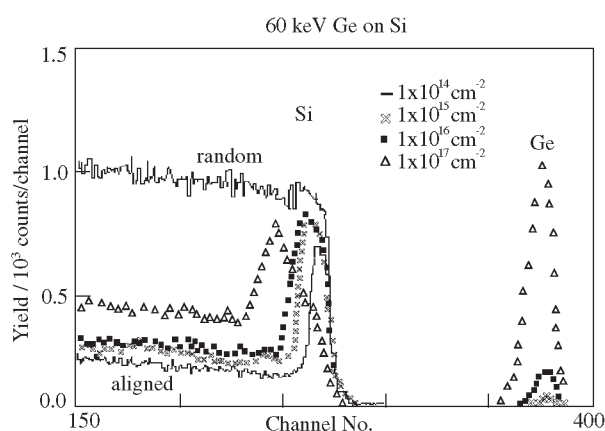


Figure 2. RBS-spectrum of a Si-sample, implanted with 60 keV Ge⁺-ions.

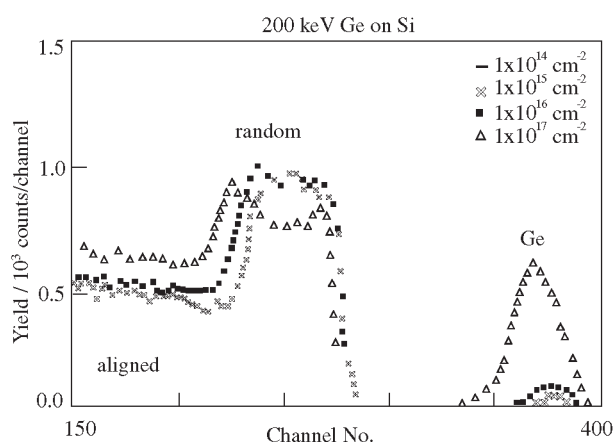


Figure 3. RBS-spectrum of a Si-sample, implanted with 200 keV Ge⁺-ions.

250 and 270. In this case the distribution of the damage resembles qualitatively the distribution of the implanted Ge⁺-ions as calculated by the Monte Carlo simulation, shown in Fig. 1, and also derived from the RBS-plots of Figs. 2 and 3, as shown in Fig. 4.

With increasing implantation dose, the intensity of this Ge backscattering peak increases and broadens, suggesting that the depth of the damaged layer as well as the extent of damage increases. Such a behavior is very likely to be due to the recoil of the Si target atoms during implantation. In addition the shift of the Si peak detected for the highest doses implies that the Ge concentration close to the surface is significantly higher than for the other samples. This concurs with the high intensity of the Ge peak, which was detected in the RBS spectra for a dose of $1 \times 10^{17} \text{ cm}^{-2}$ at channel numbers 360 to 390. The width of this Ge peak reflects the depth of the Ge distribution. The deviation from the calculated distribution of Fig. 1 is probably caused by the dif-

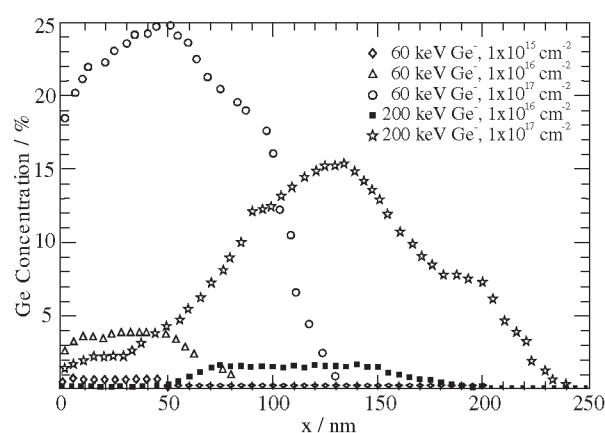


Figure 4. Ge concentration with respect to Si, derived from the RBS spectra shown in previous figures 2 and 3.

fusion of Ge atoms into the radiation damaged region. As a result the depth of the Ge distribution evaluated from the width of the Ge peak is almost identical to the depth of the radiation damaged region as measured by the width of the Si peak. This is supported by the Ge distribution derived from the RBS spectra in Fig. 4. Considering now the implantations at 200 keV ion energy, the RBS spectra of the ‘as implanted’ samples reveal a similar behavior. For the lowest dose the distribution of the damage is again in agreement with the calculated Ge distribution of Fig. 1. For the higher dose the width of the damaged layer is found to exceed the width of the Ge distribution shown in Fig. 4. Consequently the Ge ions are supposed to be impaired of diffusing into the entire damaged region. As could be deduced from the fact, that the aligned yield reach the random value (Fig. 3) for the higher implantation doses, an amorphous layer seems to be formed. This is not in contrast with the fact that the He⁻-probe ions originating from deeper regions show still channeling behavior²⁵.

Pronounced differences are manifest in the RBS spectra of Figs. 5 and 6, representing the samples which were treated by thermal annealing.

In the aligned spectra, the Si-peak denoting the radiation damage has now entirely disappeared in all cases. Obviously the crystalline structure was restored by the annealing process. The crystalline quality of the Ge-Si alloy formed can be deduced from the c_{min} of the Ge peak. At a dose of $1 \times 10^{17} \text{ cm}^{-2}$, for instance, values of 17% and 22% were found for 60 and 200 keV, respectively. Furthermore, the line shape of this peak at 60 keV also indicates that the interface between the Ge-Si alloy and the underlying Si substrate is very abrupt. This behavior is less pronounced for the RBS spectra given in Fig. 6.

From the present data it seems that the diffusion of the

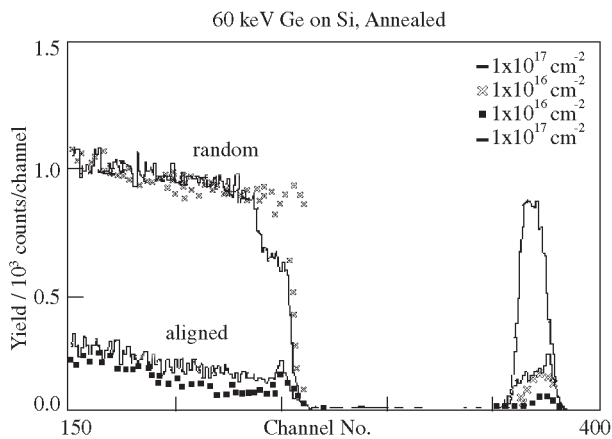


Figure 5. RBS-spectrum of a Si-sample, implanted by 60 keV Ge⁺-ions and annealed at 900 °C for 60 minutes.

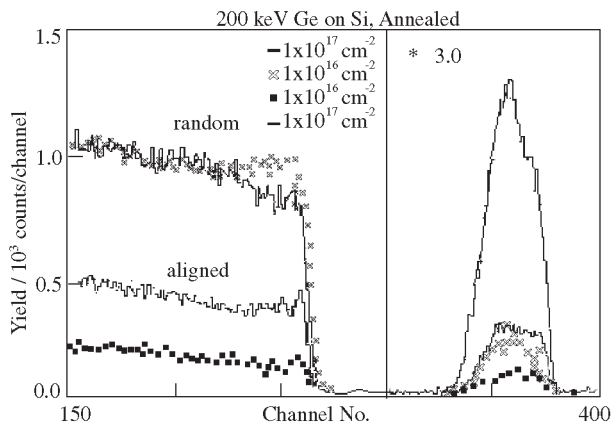


Figure 6. RBS-spectrum of a Si-sample, implanted by 200 keV Ge⁺-ions and annealed at 900 °C for 60 minutes.

Ge atoms is essentially limited to the radiation-damaged region. Consequently homogeneous Ge distributions are achievable by annealing, which is clearly demonstrated in Fig. 7 by the data of the 60 keV samples.

4. Results and Discussion

A single-crystal quality Ge-Si alloy can be produced in Si substrates by Ge⁺-ion implantation and a subsequent thermal annealing process. We found a homogeneous distribution of Ge in the layer between the Si surface and the maximum penetration depth of implanted ions is. The stoichiometric composition of the layer depends on the ion dose. By use of a fine-focus Ge⁺-ion beam, the formation of very small patterns within a Si substrate would be possible, turning this approach potentially to a useful technique for three-dimensional microelectronic device integration. The under-

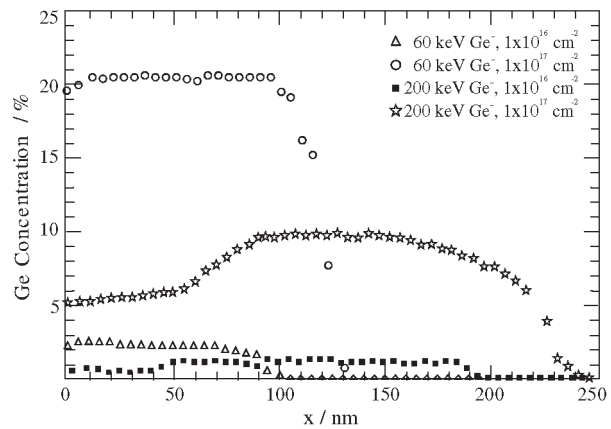


Figure 7. Germanium distribution derived from the RBS spectra shown in previous Figs. 5 and 6.

standing of key issues, such as diffusion kinetics of dopants in the Si-Ge region, device related material properties, such as carrier mobility and life time in implantation-formed Si-Ge are a future goal for the successful fabrication of ion implanted Si-Ge heterojunction transistors.

Acknowledgment

A.T. thanks Dr. Zahn, Berlin, R. Grötzchel and R. Klages, Dresden-Rossendorf, for assistance in sample preparation and measurements. A.R. is indebted to CONACyT for providing financial support through contract number 89628.

References

1. Umapathi, B.; Das, S.; Lahiri, S.K.; Kal, S. *J. Electronic Materials*, v. 30, p. 17-22, 2001.
2. Kwok, K.H. *Solid-State Electronics*, v. 43, n. 2, p. 275-283, 1999.
3. DiCaspere, L.; Capellini, L.; Cianci, E.; Evangelisti, F. *J. Vac. Sci. Technol.*, v. B16(3), p. 1721, 1998.
4. Gluck, M.; Hackbarth, T.; König, V.; Haas, A.; Hock, G.; Kohn, E. *J. Electron. Lett.*, v. 33, p. 325, 1997.
5. Mukarami, E.; Nakagawa, K.; Nishida A.; Miyao, M. *IEEE Trans. Electron. Devices*, v. 41, p. 857, 1994.
6. Selvakumar, R.C.; Hecht, B. *EDL*, v. 12, n. 8, p. 444-446, 1991.
7. Schüppen, A.; Erben, U.; Gruhle, A.; Kibbel, H.; Schumacher H.; König, U. *Proc. IEEE, IEDM 95 Tech. Dig.*, p. 743, 1995.
8. Kasper, E.; Kibbel, H.; Herzog, H.; Gruhle, A. *J. Appl. Phys.*, v. 33, p. 2415, 1994.
9. Zhou, G.L.; Morkog, H. *Thin Solid Films*, v. 231, p. 125, 1993.

10. Wang, W.C.; Denton, J.P.; Neudeck, G.W.; Ming Lee, I.; Takoudis, C.G.; Koh, M.T.K.; Kvam, E.P. *J. Vac. Sci. Technol.*, v. B15(1), p. 138, 1997.
11. Bean, J.C. *J. Cryst. Growth*, v. 81, p. 411, 1987.
12. Zehe, A. in: *Crystal Growth and Charact. Adv. Materials* (A. N. Christensen ed.), World Scientific Publ., Singapore 1987.
13. Cheung, W.Y.; Wong, S.P.; Wilson, I.H.; Zhang, T.; Chu, P.K. in: *Beam-Solid Interactions for Materials Synthesis and Characterization*, (eds. D. E. Luzzi, T. F. Heinz, M. Iwaki & D. C. Jacobson), Pittsburgh, PA, MRS Proc 354 (1995) 201-206.
14. Greve, D.W.; Raganelli, M. *J. Vac. Sci. Technol.*, v. B8, p. 511, 1990.
15. Gronet, C.M.; King, C.A.; Opyd, W.; Gibbons, J.F.; Wilson, S.D.; Hull, R. *J. Appl. Phys.*, v. 61, p. 2407, 1987.
16. De Boer, W.B.; Meyer, D.J. *Appl. Phys. Lett.*, v. 58, p.286, 1991.
17. Jang, S.M.; Tsai, C.; Reif, R. *J. Electron. Mater.*, v. 20, p. 91, 1991.
18. Abelson, J.R.; Sigmon, T.W.; Kim, K.; Weiner, K.H.; *Appl. Phys. Lett.*, v. 52(3), p. 230, 1988.
19. Wada, T.; Takeda, M.; Yamada, T.; Ichimura, M. *Appl. Surf. Sci.*, v. 41/42, p. 580, 1989.
20. Merzey, G.; Matteson, S.M.; Gyulai, J. *Nucl. Instr. and Meth.*, v. 182/183, p. 587, 1981.
21. Cheung, W.Y.; Wong, S.P.; Wilson, I.H.; Zhang, T. *Nucl. Instrum. & Meth.*, v. B 101, p. 243-246, 1995.
22. Berti, M.; Mazzi, G.; Drigo, A.V.; Migliori, A.; Jannitti, E.; Nicoletti, S. *Appl. Surf. Sci.*, v. 43, p. 158, 1989.
23. Kwok, K.H.; Selvakumar, R.C. *Proc. IEEE 1998 Canadian Conference on Electrical and Computer Engineering*, Waterloo, Canada, p. 878-881, May 1998.
24. Ziegler, J.F.; Biersack, J.P.; Littmark, U. *The Stopping Power and Ranges of Ions in Matter*, Vol. I, Pergamon Press 1985.
25. Lu, X.; Cheung, N.W. *Appl. Phys. Lett.*, v. 69, p. 1915, 1996.
26. Kwok, K.H.; Selvakumar, R.C. *Proc. of the 1st Waterloo Workshop on Materials Technology. Materials Technology '98*, Waterloo, Canada, 13 June 1998.